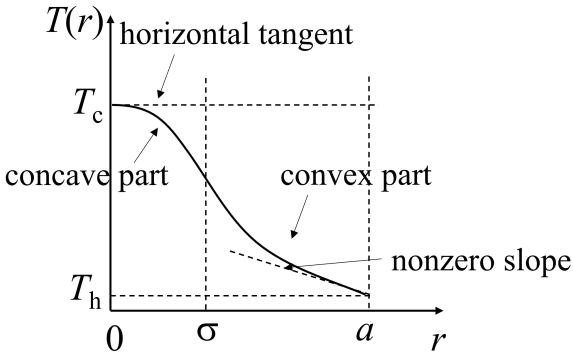


T1: Thermal lens - Solution

(a) Drawing a $T(r)$ graph

The graph should present or clearly infer for the four elements shown in the figure.



Element of the graph	Pts.
Horizontal tangent at $r = 0$	0.5
Graph is concave and decreasing in $0 \leq r \leq \sigma$	0.5
Graph is convex and decreasing in $\sigma \leq r \leq a$	0.5
Nonzero negative slope at $r = a$	0.5
Totally on (a)	2.0

(b) Finding T_c

Approach with a direct solution of the heat-transport equation.

Consider a cylindrical cut of the disk with radius r . Let $P_{\text{abs}}(r)$ be the portion of laser power absorbed within the cylinder. The absorbed power is being transferred as heat towards the outer holder through the circumvent surface $2\pi br$ of the cylinder. The heat-transport equation reads:

$$-k(2\pi rb) \frac{dT(r)}{dr} = P_{\text{abs}}(r) \quad (1)$$

Depending on r , the absorbed power is given by two different expressions. Within the illuminated area $0 \leq r \leq \sigma$, the incident light intensity $I = P_L/(\pi\sigma^2)$ is constant, so the absorbed power is:

$$P_{\text{abs}}(r) = I\pi r^2 = AP_L \frac{r^2}{\sigma^2} \quad (2)$$

and the solution for $T(r)$ is quadratic in r :

$$T(r) = T_c - \frac{AP_L r^2}{4\pi kb\sigma^2} \quad (3)$$

where $T_c = T(0)$ is the temperature at the center. It is clear from that expression that the parameter m is:

$$m = -\frac{AP_L}{4\pi kb\sigma^2} = -1.1 \cdot 10^7 \text{ Km}^{-2} \quad (4)$$

Outside the illuminated area, i.e. for $\sigma \leq r \leq a$, $P_{\text{abs}}(r) = AP_L$, which does not depend on r . The solution for $T(r)$ is logarithmic in r :

$$T(r) = T_h + \frac{AP_L}{2\pi kb} \ln\left(\frac{a}{r}\right) \quad (5)$$

where $T_h = T(a)$ is the temperature of the holder, which is equal to the temperature along the outer

rim of the disk. After matching the two solutions at $r = \sigma$ we obtain:

$$T_c = T_h + \frac{AP_L}{4\pi kb} \left[1 + 2 \ln \left(\frac{a}{\sigma} \right) \right] = 41^\circ\text{C} \quad (6)$$

Task	Pts.
Writes down the heat-transport equation in radial coordinates	0.7
Derives an expression for P_{abs} at $0 \leq r \leq \sigma$	0.5
Finds a quadratic solution for $T(r)$ in the region $0 \leq r \leq \sigma$ (Subtract 0.2 pts. if the boundary condition $T(0) = T_c$ has not been accounted for)	0.5
Identifies the expression for m (Subtract 0.1 pts. if the minus sign is missing)	0.2
Calculates m numerically (Subtract 0.1 pts. if the minus sign is missing)	0.2
Derives an expression for P_{abs} at $\sigma \leq r \leq a$	0.3
Finds a logarithmic solution for $T(r)$ in the region $\sigma \leq r < a$ (Subtract 0.2 pts. if the boundary condition $T(a) = T_h$ has not been accounted for)	0.5
Sets up an equation by matching the two solutions at $r = a$	0.6
Writes down the final expression for T_c	0.2
Calculates T_c numerically	0.3
Totally on (b)	4.0

Alternative approach - direct piece-wise integration of the heat-transport equation.

After realizing that $P_{\text{abs}}(r)$ is given by a piece-wise function:

$$P_{\text{abs}} = \begin{cases} AP_L r^2 / \sigma^2 & \text{if } 0 \leq r \leq \sigma \\ AP_L = \text{const} & \text{if } \sigma \leq r \leq a \end{cases} \quad (7)$$

the student may substitute the given solution $T(r) = T_c + mr^2$ into heat-transport equation (1) for $0 \leq r \leq \sigma$. This gives directly the expression (4) for the parameter m . The parameter T_c could be easily identified with the temperature $T(0)$ at the center of the disk. On the other hand, $T_h = T(a)$ due to the thermal contact between the rim of the disk and the holder. It follows from the heat-transport equation that:

$$-\frac{dT(r)}{dr} = \frac{P_{\text{abs}}(r)}{2\pi kbr} \quad (8)$$

The piece-wise integration of the two sides of the equation in the interval $0 \leq r \leq a$ gives:

$$\begin{aligned} T(0) - T(a) &= T_c - T_h = \int_0^a \frac{P_{\text{abs}}(r)}{2\pi kbr} dr \\ &= \int_0^\sigma \frac{P_{\text{abs}}(r)}{2\pi kbr} dr + \int_\sigma^a \frac{P_{\text{abs}}(r)}{2\pi kbr} dr \\ &= \frac{AP_L}{4\pi kb} \left[1 + 2 \ln \left(\frac{a}{\sigma} \right) \right] \end{aligned} \quad (9)$$

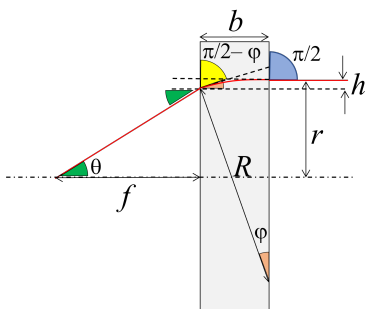
which is equivalent to the expression (6) for T_c .

Task	Pts.
Writes down the heat-transport equation in radial coordinates	0.7
Derives an expression for P_{abs} at $0 \leq r \leq \sigma$	0.5
Derives an expression for P_{abs} at $\sigma \leq r \leq a$	0.3
Substitutes the given form of $T(r)$ into heat-transport equation for the region $0 \leq r \leq \sigma$	0.3
Obtains the expression for m (Subtract 0.1 pts. if the minus sign is missing)	0.2
Calculates m numerically (Subtract 0.1 pts. if the minus sign is missing)	0.2
States that $T_c = T(0)$	0.2
States that $T_h = T(a)$	0.2
Expresses $T_c - T_h$ through integral of dT/dr in the $0 \leq r \leq a$ interval	0.3
Calculates the integral in the $0 \leq r \leq \sigma$ interval	0.3
Calculates the integral in the $\sigma \leq r \leq a$ interval	0.3
Writes down the final expression for T_c	0.2
Calculates T_c numerically	0.3
Totally on (b)	4.0

(c) Finding the focal length

Approach based on the Fermat's principle

We consider only the illuminated area of the disk ($0 \leq r \leq \sigma$). Due to the nonuniform temperature distribution, the index of refraction is also r -dependent, which leads to a bending of the light rays incident at nonzero radii r , as shown schematically in the figure. As a result, the light rays exiting the disk, converge toward the optical axis, and, eventually, cross it in a certain point at a distance f from the disk.



Since the ray bending is relatively small, one may assume that: (i) all the rays travel approximately the same distance b inside the disk; (ii) any given ray enters and exits the disk at approximately the same height r ; (iii) $f \gg r$. Thus, the optical pathway $s(r)$ of a ray, incident at a height r above the optical axis, is:

$$s(r) \approx n(r)b + \sqrt{f^2 + r^2} \approx n(r)b + f + \frac{r^2}{2f} \quad (10)$$

According to Fermat's principle, in order that all the rays focus at the same point, it is necessary that $s(r)$ is constant within $0 \leq r < \sigma$. In particular, $s(r) \equiv s(0)$ for any r , which leads to the condition:

$$b(n(0) - n(r)) \equiv \frac{r^2}{2f} \quad (11)$$

Since:

$$n(0) - n(r) = \gamma(T(0) - T(r)) = -\gamma m r^2 = \gamma |m| r^2 \quad (12)$$

condition (11) is satisfied if

$$\gamma b|m|r^2 \equiv \frac{r^2}{2f} \quad (13)$$

The beam hence focuses at

$$f = \frac{1}{2\gamma b|m|} \quad (14)$$

Taking into account the expression for m derived in part (b), we represent the answer in terms of the known parameters and calculate its numerical value:

$$f = \frac{2\pi k\sigma^2}{\gamma AP} \approx 0.94 \text{ m} \quad (15)$$

Alternatively, the student may state that the optical pathway $s(r)$ does not depend on r and use the condition $ds/dr \equiv 0$, which gives:

$$\frac{dn(r)}{dr}b + \frac{r}{f} \equiv 0 \quad (16)$$

Since:

$$\frac{dn(r)}{dr} = \frac{dn}{dT} \frac{dT}{dr} = \gamma 2mr \quad (17)$$

we obtain:

$$\frac{r}{f} + 2\gamma bmr \equiv 0 \quad (18)$$

The beam thus focuses at:

$$f = \frac{1}{2\gamma b|m|} = \frac{2\pi k\sigma^2}{\gamma AP} \approx 0.94 \text{ m} \quad (19)$$

Task	Pts.
Formulates assumptions (i) and (ii) or equivalent statements	0.2
Derives a general expression for the optical pathway s as a function of r	0.8
Uses that $f \gg r$ and derives approximate quadratic expression for $s(r)$	0.5
States that focusing takes place when $s(r)$ is the same for all rays converging in the focus	0.5
Writes explicitly equation in the form $s(r) \equiv s(0)$ OR $ds/dr \equiv 0$	0.5
Uses the $T(r)$ dependence to derive the $n(r)$ dependence OR to find dn/dr	0.5
Derives an expression for the focal length f in terms of m or of the parameters given in the problem statement	0.8
Calculates f numerically	0.2
Totally on (c)	4.0

Approach based on a direct ray/wavefront tracing

As a next approximation, the light ray inside the disk can be modeled as a circular arc of a radius R ($R \gg b$, see the figure). As a result, the ray exits the disk at a smaller height $r - h$ above the optical axis ($h \ll r$). The angle of bending φ of the ray inside the material is related to h by:

$$\cos \varphi = 1 - \frac{h}{R} \quad (20)$$

From the Snell's law it follows that:

$$n(r) \sin(\pi/2) = n(r-h) \sin(\pi/2 - \varphi) = n(r-h) \cos \varphi \quad (21)$$

Up to terms, linear in h , one may write that:

$$\cos \varphi = \frac{n(r)}{n(r-h)} \approx 1 + \frac{n'(r)}{n(r)} h \quad (22)$$

Thus, the radius of the ray inside the material, is:

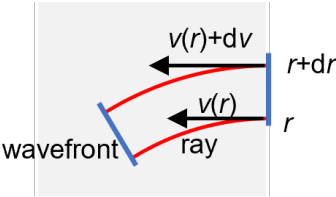
$$R = -\frac{n(r)}{n'(r)} \quad (23)$$

and the bending angle is approximately:

$$\varphi = \frac{b}{R} = -\frac{n'(r)b}{n(r)} \quad (24)$$

Alternatively, the students may trace a small part of the wavefront, associated with two rays, incident at close distances r and $r + dr$ from the optical axis. By noticing that the wavefront is perpendicular to the rays, the angle φ of deviation of the rays is equal to the angle of rotation of the wavefront. Let $v(r) = c/n(r)$ be the speed of light in the material at a distance r from the axis. The time-rate $\dot{\varphi}$, i.e. the angular speed of the wavefront is:

$$\dot{\varphi} = dv(r)/dr = -cn'(r)/n(r)^2 \quad (25)$$



The rays reach the opposite surface of the disk in approximately the same time $t = bn(r)/c$, so the total deflection angle of the wavefront, and of the rays, thereof, is $\varphi = \dot{\varphi}t = -n'(r)b/n(r)$

Upon exiting the disk, the ray undergoes additional refraction, and inclines at a new angle θ relative to the optical axis. The Snell's law in the small-angle approximation ($\sin \theta \approx \theta$, $\sin \varphi \approx \varphi$) states that:

$$\theta = n(r-h)\varphi \approx n(r)\varphi = -n'(r)b \quad (26)$$

Since $n'(r) = \gamma T'(r) = 2mr$, one obtains the following expression for the angle of inclination:

$$\theta = 2\gamma|m|r \quad (27)$$

It is clear from the figure that:

$$f = \frac{r-h}{\tan \theta} \approx \frac{r}{\theta} = \frac{1}{2\gamma|m|b} \quad (28)$$

which reproduces the result obtained by the Fermat's principle.

Task	Pts.
States or shows on a graph that the light ray inside the material can be approximated by an arc OR illustrates on a graph that the wavefront is perpendicular to the light rays	0.2
Derives the relation $\cos \varphi = 1 - h/R$ or equivalent OR expresses the time of travel of the ray across the disk	0.5
Applies the Snell's law to the ray path inside the material OR derives a formula for the time-rate $\dot{\varphi}$ by considering the wavefront passing through two closely separated rays	0.8
Finds expression for the bending angle φ inside the material	0.5
Applies the Snell's law for the refraction of rays exiting the disk	0.5
Uses the $T(r)$ dependence to find $n'(r)$	0.5
Derives an expression for the focal length f in terms of m or of the parameters given in the problem statement	0.8
Calculates f numerically	0.2
Totally on (c)	4.0

Alternatively:

Consider a ray incident at a specific height r_0 . Let $x \in [0, b]$ be the horizontal coordinate of the ray inside the material. Our goal is to find approximately the profile $r(x)$ of the light ray inside the material. From Snell's law:

$$n(r) \cos \varphi = n(r_0)$$

Since:

$$\tan \varphi = -\frac{dr(x)}{dx} = -\frac{\sqrt{(1 - \cos^2 \varphi)}}{\cos \varphi}$$

we obtain:

$$-\frac{dr(x)}{dx} = \frac{\sqrt{(n(r) - n(r_0))(n(r) + n(r_0))}}{n(r)}$$

Since $n(r) \approx n(r_0) + n'(r_0)(r - r_0) = n(r_0) + 2\gamma|m|r_0(r_0 - r)$, the differential equation approximates to:

$$-\frac{dr(x)}{dx} = \sqrt{\frac{4\gamma|m|r_0}{n(r_0)}} \sqrt{r_0 - r}$$

which is solvable by separation of variables and gives an approximate parabolic path:

$$r(x) = r_0 - \frac{\gamma|m|r_0}{n(r_0)} x^2$$

Finally, for the angle, just before exiting the disk, we obtain:

$$\varphi \approx \tan \varphi = -r'(x = b) = \frac{2\gamma|m|r_0 b}{n(r_0)}$$

Task	Pts.
Applies the Snell's law to express the angle φ dependence on r for one ray	0.8
Expresses $\tan \varphi = dr/dx$	0.2
Applies approximation to derive a separable differential equation for $r(x)$	0.8
Integrates to obtain parabolic path	0.5
Finds expression for the angle φ , before exiting the disk	0.2
Applies the Snell's law for the refraction of rays exiting the disk	0.5
Derives an expression for the focal length f in terms of m or of the parameters given in the problem statement	0.8
Calculates f numerically	0.2
Totally on (c)	4.0